

T_0 Separation in Axiomatic Quantum Mechanics[†]

Bart Van Steirteghem^{1,2}

Received December 8, 1999

Using the physical duality between states and properties, Aerts *et al.* obtained a “lattice” representation for all closure spaces, through state property systems. In this paper I discuss the equivalence of ‘state determination’ for state property systems with T_0 separation for closure spaces. I also provide a link with well-known lattice representations of closure spaces, through some results of Ern e.

1. DUALITY OF STATES AND PROPERTIES

In the Geneva–Brussels approach to the foundations of physics (Piron, 1976, 1989, 1990; Aerts, 1982, 1983; Moore, 1999) an entity is described by its ‘states’ and its ‘properties.’ The physical relation between states and properties is ‘actuality.’ So, I presuppose the existence of a relation I between the set \mathcal{L} of properties and the set Σ of states,³ where for $a \in \mathcal{L}$ and $p \in \Sigma$, aIp is interpreted as “property a is actual in state p .” As with any relation, I induces a closure on Σ and on \mathcal{L} . This is a result due to Birkhoff (1967, Chapter 5, §7). Indeed, if for $A \subset \Sigma$ and $B \subset \mathcal{L}$ one puts

$$A^+ = \{a \in \mathcal{L} | aIp \ \forall p \in A\} \quad (1)$$

$$B^* = \{p \in \Sigma | aIp \ \forall a \in B\} \quad (2)$$

then $\mathcal{P}(\Sigma) \rightarrow \mathcal{P}(\Sigma)$, $A \mapsto A^{+*}$, and $\mathcal{P}(\mathcal{L}) \rightarrow \mathcal{P}(\mathcal{L})$, $B \mapsto B^{*\dagger}$, are closure operators. Moreover, Birkhoff proves that the mappings $A \mapsto A^+$ and $B \mapsto B^*$ define dual isomorphisms between the complete lattices of closed subsets of Σ and \mathcal{L} .

[†]This paper is dedicated to the memory of Prof. Gottfried T. R uttimann.

¹FUND, Department of Mathematics, Brussels Free University, B-1050 Brussels, Belgium; E-mail: bvsteirt@vub.ac.be.

²Research Assistant, Fund for Scientific Research—Flanders.

³In the language of formal concept analysis (Ganter and Wille, 1999), (\mathcal{L}, Σ, I) is a ‘context.’

The relation I encodes the “physical duality of states and properties” (Moore, 1999). Using I , one can define a preorder on \mathcal{L} :

$$a < b \Leftrightarrow a^* \subset b^* \quad (3)$$

with a slight abuse of notation. It is of course “physical implication with respect to actuality” (Moore, 1999). After identification of mutually implying properties, \mathcal{L} becomes a poset. It is then operationally justified (a key result in the Geneva School) to demand that \mathcal{L} is a complete lattice in which the meet is “physical conjunction with respect to actuality” (Moore, 1999):

$$(\wedge_i a_i)Ip \Leftrightarrow a_iIp \quad \forall_i \quad (4)$$

which is read as “ $\wedge_i a_i$ is actual iff every a_i is actual.” This is equivalent to demanding every closed subset of Σ be of the form a^* for some property $a \in \mathcal{L}$. Indeed, let $F = F^{\dagger*}$ be a closed subset of Σ . From (4) it follows that $(\wedge F^\dagger)^* = \bigcap_{a \in F^\dagger} a^* = F^{\dagger*} = F$. Conversely, consider a family $a_i \in \mathcal{L}$. Then there exists $b \in \mathcal{L}$ such that $b^* = \bigcap_i a_i^*$. Consequently, $b = \wedge_i a_i$ and (4) holds. In the Geneva School, the mapping $a \mapsto a^*$, which is now a lattice isomorphism between \mathcal{L} and the lattice of closed subsets of Σ , has been named the *Cartan map*.

Putting these ideas into one structure and writing $\xi(p)$ for p^\dagger ($p \in \Sigma$), Aerts *et al.* (1999) defined state property systems. A triple $(\Sigma, \mathcal{L}, \xi)$ is a *state property system (sps)* if Σ is a set, $(\mathcal{L}, \wedge, <)$ is a complete lattice, and $\xi: \Sigma \rightarrow \mathcal{P}(\mathcal{L})$ is a function such that $\xi(p)$ never contains the universal lower bound 0 of \mathcal{L} (0 is never actual) and

$$a < b \Leftrightarrow \text{if } a \in \xi(p) \text{ then } b \in \xi(p) \quad (5)$$

$$\wedge_i a_i \in \xi(p) \Leftrightarrow a_i \in \xi(p) \quad \forall_i \quad (6)$$

Obviously (5) and (6) are restatements of (3) and (4). Putting $s_\xi(p) = \wedge \xi(p)$, it is clear that $\xi(p) = [s_\xi(p), 1]$ for every $p \in \Sigma$ (1 is the maximum of \mathcal{L}). Evidently, $s_\xi(p)$ is the strongest (minimal) property which is actual in state p .

A couple (m, n) is a morphism $(\Sigma', \mathcal{L}', \xi') \rightarrow (\Sigma, \mathcal{L}, \xi)$ of sps's if $m: \Sigma' \rightarrow \Sigma$ and $n: \mathcal{L}' \rightarrow \mathcal{L}$ are maps such that for $p' \in \Sigma'$ and $a \in \mathcal{L}$

$$a \in \xi(m(p')) \Leftrightarrow n(a) \in \xi'(p') \quad (7)$$

For the physical idea behind this definition, I refer to Aerts *et al.* (1999). The category of sps's and their morphisms is denoted **SP**. The two functors given in (8) and (9) below establish an equivalence between **SP** and the category of closure spaces with \emptyset closed and continuous maps **Cls** (Aerts *et al.*, 1999).

$$F: \mathbf{SP} \rightarrow \mathbf{Cls}, \quad \begin{cases} (\Sigma, \mathcal{L}, \xi) \mapsto (\Sigma, \mathcal{F}_{\mathcal{L}}) \\ (m, n) \mapsto m \end{cases} \quad (8)$$

$$G: \mathbf{Cls} \rightarrow \mathbf{SP}, \quad \begin{cases} (\Sigma, \mathcal{F}) \mapsto (\Sigma, \overline{\mathcal{F}}, \xi_{\overline{\mathcal{F}}}) \\ m \mapsto (m, m^{-1}) \end{cases} \quad (9)$$

where $\mathcal{F}_{\mathcal{L}} = \{\{p \in \Sigma: a \in \xi(p)\} \mid a \in \mathcal{L}\}$ and $\xi_{\overline{\mathcal{F}}}: \Sigma \rightarrow \mathcal{P}(\overline{\mathcal{F}})$, $p \mapsto \{F \in \overline{\mathcal{F}} \mid p \in F\}$. So, on the object side, F constructs one of the “Birkhoff polarity closures” for $I \subset \mathcal{L} \times \Sigma$ defined by aIp if $a \in \xi(p)$. Conversely, the object correspondence of G is presented in Aumann (1970). This equivalence provides a “lattice representation” for all closure spaces and, in this sense, generalizes Ern e’s (1984) lattice representation for T₀ closure spaces. As will be explained in Section 4, the latter representation is the cornerstone of Ern e’s general construction providing “all” lattice representable categories of closure spaces. In the same section I shall show how the above equivalence fits into this scheme.

In the final section I give a short discussion of T₀ separation for orthogonality spaces.

2. STATE DETERMINATION AND T₀ SEPARATION

Traditionally, in the Geneva–Brussels approach, the state p of an entity is identified with its actual properties, i.e., with $\xi(p)$. In this section I review the implications of this (physical) assumption on the equivalence of **SP** and **Cls**. Since they are straightforward, I omit the proofs, which can be found in Aerts *et al.* (1999).

We call an sps $(\Sigma, \mathcal{L}, \xi)$ *state determined* if ξ is injective.⁴ In other words, $(\Sigma, \mathcal{L}, \xi)$ is state determined if every state $p \in \Sigma$ is determined by its actual properties $\xi(p)$. Let **SP**₀ be the full subcategory of **SP**, with state-determined sps’s as objects.

Lemma 1. Let $(\Sigma, \mathcal{L}, \xi)$ be an sps. The following are equivalent:

1. $(\Sigma, \mathcal{L}, \xi)$ is state determined.
2. $s_{\xi}: \Sigma \rightarrow \mathcal{L}, p \mapsto \wedge \xi(p)$, is injective.
3. $F(\Sigma, \mathcal{L}, \xi)$ is a T₀ closure space.

Conversely, a closure space (Σ, \mathcal{F}) is T₀ iff $G(\Sigma, \mathcal{F}) = (\Sigma, \overline{\mathcal{F}}, \xi_{\overline{\mathcal{F}}})$ is a state-determined sps.

Recall that a closure space (Σ, \mathcal{F}) is said to be T₀ if $cl(p) = cl(q) \Rightarrow p = q$ for $p, q \in \Sigma$, where for $A \subset \Sigma$, $cl(A) \doteq \bigcap \{F \in \mathcal{F} \mid A \subset F\}$.

⁴In formal concept analysis the corresponding contexts are called ‘clarified.’

Let \mathbf{Cls}_0 be the full subcategory of \mathbf{Cls} of T_0 closure spaces. The next proposition easily follows.

Proposition 1. The functors F and G of (8) and (9) restrict and corestrict to equivalence establishing functors between \mathbf{SP}_0 and \mathbf{Cls}_0 .

3. STATES AS STRONGEST ACTUAL PROPERTIES

Let $(\Sigma, \mathcal{L}, \xi)$ be a state-determined sps. Then, by condition 2 of Lemma 1, a state $p \in \Sigma$ may be identified with $s_\xi(p) \in \mathcal{L}$, i.e., with the strongest property it makes actual. As a consequence, Σ can be embedded into \mathcal{L} as an order-generating subset (see Lemma 2). This engenders another equivalence of categories (Proposition 2). The proofs can again be found in Aerts *et al.* (1999).

Lemma 2. Let $(\Sigma, \mathcal{L}, \xi)$ be an sps. Then $0 \notin \Sigma^\xi \doteq s_\xi(\Sigma)$ is an order-generating subset of \mathcal{L} : for every a in \mathcal{L} ,

$$a = \vee \{x \in \Sigma^\xi \mid x < a\} \tag{10}$$

A couple (Σ, \mathcal{L}) is a *based complete lattice (bcl)* if \mathcal{L} is a complete lattice and $\Sigma \subset \mathcal{L}$ is an order-generating subset not containing 0. The previous lemma then becomes $(\Sigma, \mathcal{L}, \xi) \in \mathbf{SP} \Rightarrow (\Sigma^\xi, \mathcal{L})$ is a bcl.

Lemma 3. Let (Σ, \mathcal{L}) be a bcl. If we define

$$\xi: \Sigma \rightarrow \mathcal{P}(\mathcal{L}), p \mapsto [p, 1] \tag{11}$$

then $(\Sigma, \mathcal{L}, \xi)$ is a state-determined sps.

To deal with the morphisms, I shall use Galois connections. I shall write n_* for the lower adjoint of a meet-preserving map n and f^* for the upper adjoint of a join-preserving f .

Consider two bcl's (Σ', \mathcal{L}') , (Σ, \mathcal{L}) . A function $\bar{f}: \mathcal{L}' \rightarrow \mathcal{L}$ is a *morphism of bcl's* if f preserves joins and $f(\Sigma') \subset \Sigma$. The category of bcl's will be denoted \mathbf{L}_0 .

Proposition 2. The following two functors establish an equivalence between \mathbf{SP}_0 and \mathbf{L}_0 :

$$H: \mathbf{SP}_0 \rightarrow \mathbf{L}_0, \quad \begin{cases} (\Sigma, \mathcal{L}, \xi) \mapsto (\Sigma^\xi, \mathcal{L}) \\ (m, n) \mapsto n_* \end{cases} \tag{12}$$

$$K: \mathbf{L}_0 \rightarrow \mathbf{SP}_0, \quad \begin{cases} (\Sigma, \mathcal{L}) \mapsto (\Sigma, \mathcal{L}, \xi) \\ f \mapsto (f|_{\Sigma'}, f^*) \end{cases} \tag{13}$$

where ξ of (13) is given in (11).

The results above can be summarized in the following scheme, which can be read as a commutative diagram.

$$\begin{array}{ccc} \mathbf{Cls} & \approx & \mathbf{SP} \\ \cup & & \cup \\ \mathbf{Cls}_0 & \approx & \mathbf{SP}_0 \approx \mathbf{L}_0 \end{array} \quad (14)$$

4. CONNECTION WITH ERNÉ’S RESULT

Erné (1984) gives a direct proof of the equivalence of **Cls**₀ and **L**₀. In fact, his functors are $H \circ G$ and $F \circ K$. Based on this equivalence, he gives a general construction for “lattice representable” (or *l-representable*, after Ern e) categories of closure spaces. I give an outline of his result. Note that I shall change his definition of ‘invariant selection’ slightly: for Ern e the empty set need not be closed.

Given a category **C** of closure spaces, i.e., a full and isoclosed subcategory of **Cls**, Ern e introduces the *l*-representing functor

$$T: \mathbf{C} \rightarrow \mathbf{L}_\vee, \quad \begin{cases} (\Sigma, \mathcal{F}) \mapsto (\mathcal{F}, \cap, \subset) \\ f \mapsto [F \mapsto cl(f(F))] \end{cases} \quad (15)$$

where **L**_∨ is the category of complete lattices with join-preserving maps. **C** is then called *l-representable* if a suitable corestriction of T yields an equivalence of categories. Ern e gives many examples of such categories. Well known is the equivalence of the category of T_1 closure spaces and the category of complete atomistic lattices.

Next, let **L** be an isoclosed subclass of the class of complete lattices. An *invariant selection* S for **L** is a (class-theoretic) function assigning to each $\mathcal{L} \in \mathbf{L}$ a certain $0 \notin S(\mathcal{L}) \subset \mathcal{L}$, such that whenever ψ is a lattice isomorphism between **L**-elements \mathcal{L} and \mathcal{L}' , then $\psi(S(\mathcal{L})) = S(\mathcal{L}') = S(\psi(\mathcal{L}))$. Given an invariant selection S , Ern e defines the isoclosed subcategory **L** _{S} of **L**_∨ as follows. A complete lattice $\mathcal{L} \in \mathbf{L}$ is an object of **L** _{S} iff $S(\mathcal{L})$ is an order-generating subset of \mathcal{L} . An **L** _{S} -morphism $\varphi: \mathcal{L} \rightarrow \mathcal{L}'$ is a join-preserving map, such that $\varphi(S(\mathcal{L})) \subset S(\mathcal{L}')$. Hence, **L** _{S} may be (and is) considered a full subcategory of **L**₀.

He calls a closure space (Σ, \mathcal{F}) *S-complete* if it is T_0 , $T(\Sigma, \mathcal{F}) \in \mathbf{L}$, and $S(T(\Sigma, \mathcal{F})) = \{cl(p) \mid p \in \Sigma\}$. The *S*-complete closure spaces form an isoclosed and full subcategory of **Cls**₀, denoted **C** _{S} . Finally, I state Ern e’s theorems.

Theorem 1. For every invariant selection S , the categories **C** _{S} and **L** _{S} are equivalent.

This equivalence is a restriction of the equivalence between **Cls**₀ and **L**₀. The next theorem says that all *l*-representable categories of closure spaces can be obtained this way.

Theorem 2. A category \mathbf{C} of closure spaces is l -representable if and only if there exists an isoclosed class \mathbb{L} of complete lattices and an invariant selection S for \mathbb{L} such that $\mathbf{C} = \mathbf{C}_S$.

Given an isoclosed class \mathbb{L} and an invariant selection S , I also introduce the full subcategory \mathbf{SP}_S of \mathbf{SP} of state property systems $(\Sigma, \mathcal{L}, \xi)$ such that $F(\Sigma, \mathcal{L}, \xi) = (\Sigma, \mathcal{F}_\xi)$ is in \mathbf{C}_S . Using that $F \circ G$ is the identical functor, it is obvious that F and G (co)restrict to an equivalence $\mathbf{C}_S \approx \mathbf{SP}_S$. Recall that the (Cartan) isomorphism κ between \mathcal{L} and \mathcal{F}_ξ is given by $\kappa(a) = \{p \in \Sigma \mid a \in \xi(p)\}$. If $(\Sigma, \mathcal{L}, \xi)$ belongs to \mathbf{SP}_S , then $\mathcal{F}_\xi \in \mathbb{L}$ and $S(\mathcal{F}_\xi) = S(\kappa(\mathcal{L})) = \{cl(p) \mid p \in \Sigma\}$. Therefore, since $\mathcal{L} \cong \mathcal{F}_\xi$, $\mathcal{L} \in \mathbb{L}$ and $S(\mathcal{L}) = \kappa^{-1}(S(\kappa(\mathcal{L}))) = \Sigma^\xi$. It follows that $H(\Sigma, \mathcal{L}, \xi)$ is in $\mathbf{L}_S \subset \mathbf{L}_0$. Using that $F \circ K: \mathbf{L}_S \rightarrow \mathbf{C}_S$ (this is one of Ern e’s functors in Theorem 1), it is now straightforward that H and K (co)restrict to an equivalence $\mathbf{SP}_S \approx \mathbf{L}_S$.

Summarizing, any l -representable category of closure spaces \mathbf{C} fits, for a suitable invariant selection S , into the following scheme, which can be read as a commutative diagram and which shows how the equivalences of Sections 1–3 refine Ern e’s beautiful result:

$$\begin{array}{ccc}
 \mathbf{Cls} & \approx & \mathbf{SP} \\
 \cup & & \cup \\
 \mathbf{Cls}_0 & \approx & \mathbf{SP}_0 \approx \mathbf{L}_0 \\
 \cup & & \cup \quad \cup \\
 \mathbf{C} = \mathbf{C}_S & \approx & \mathbf{SP}_S \approx \mathbf{L}_S
 \end{array} \tag{16}$$

5. ORTHOGONALITY SPACES AND T_0 SEPARATION

Let Σ be a set and let \perp be an antireflexive and symmetric relation on Σ . I then call (Σ, \perp) a *pseudo orthogonality space (pos)*. If for $A \subset \Sigma$, $A^\perp = \{p \in \Sigma \mid p \perp a \ \forall a \in A\}$, then $\mathcal{P}(\Sigma) \rightarrow \mathcal{P}(\Sigma)$, $A \mapsto A^{\perp\perp}$ is a closure operator such that $\emptyset^{\perp\perp} = \emptyset$. Moreover, $A \mapsto A^\perp$ is an orthocomplementation on the lattice of closed (biorthogonal) subsets (Birkhoff, 1967, Chapter 5, §7; Moore, 1995). Moore (1995) has proven that this closure is T_1 iff it separates points: $p \neq q \Rightarrow \exists r: p \perp r, q \not\perp r$; which is equivalent to $p \neq q \Rightarrow p^\perp \not\subset q^\perp$. The couple (Σ, \perp) is then called a ‘state space’ or an ‘orthogonality space’ (Moore, 1995, 2000). T_0 separation can be characterized analogously.

Proposition 3. Let \perp be a symmetric relation on Σ . The closure operator $\mathcal{P}(\Sigma) \rightarrow \mathcal{P}(\Sigma)$, $A \mapsto A^{\perp\perp}$, is T_0 iff

$$p \neq q \Rightarrow p^\perp \neq q^\perp \tag{17}$$

Indeed, if (17) holds and $p^{\perp\perp} = q^{\perp\perp}$, then $p^{\perp\perp\perp} = p^\perp = q^\perp$, whence $p =$

q . Conversely, if $p^\perp = q^\perp$, then $p^{\perp\perp} = q^{\perp\perp}$, and so by T₀, $p = q$. In fact, the mathematics behind Lemma 1 and this Proposition are the same.

Consider two pos's (Σ_1, \perp_1) and (Σ_2, \perp_2) and the following three pos's:

1. $\Sigma_1 \times \Sigma_2$ with $(p_1, p_2) \perp_{\Pi} (q_1, q_2) \Leftrightarrow p_1 \perp_1 q_1$ and $p_2 \perp_2 q_2$.
2. $\Sigma_1 \dot{\cup} \Sigma_2$ with $p_i \perp_{\Pi} q_j \Leftrightarrow i \neq j$ or $(i = j \text{ and } p_i \perp_i q_i)$ ($i, j \in \{1, 2\}$).
3. $\Sigma_1 \times \Sigma_2$ with $(p_1, p_2) \perp_{\otimes} (q_1, q_2) \Leftrightarrow p_1 \perp_1 q_1$ or $p_2 \perp_2 q_2$.

The closure associated to the first pos is the product (in **Cls**) of $(\Sigma_1, \cdot^{\perp\perp_1})$ and $(\Sigma_2, \cdot^{\perp\perp_2})$. The second pos generates the coproduct $(\Sigma_1, \cdot^{\perp\perp_1}) \amalg (\Sigma_2, \cdot^{\perp\perp_2})$. The last one is the 'separated product' introduced by Aerts (1982). For the first two it is well known that they are T₀ (T₁) iff Σ_1 and Σ_2 are. That this is also true for the separated product in the T₁ case was shown by Aerts (1982, Theorem 26).

Proposition 4. The separated product of Σ_1 and Σ_2 is T₀ iff Σ_1 and Σ_2 are.

Sufficiency can be proven as follows. Take $(p_1, p_2) \neq (q_1, q_2)$ and suppose that $p_1 \neq q_1$ and $p_1^\perp \ni r_1 \notin q_1^\perp$. Then $(p_1, p_2) \perp (r_1, q_2) \not\perp (q_1, q_2)$. For necessity one can assume Σ_1 and Σ_2 are nonempty. Consider $p_1 \neq q_1$ in Σ_1 and take $r_2 \in \Sigma_2$. Then $(p_1, r_2)^\perp \neq (q_1, r_2)^\perp$. Suppose there is a (t_1, t_2) such that $(p_1, r_2) \perp (t_1, t_2) \not\perp (q_1, r_2)$; then $p_1 \perp t_1 \not\perp q_2$.

ACKNOWLEDGMENT

I thank a referee for pointing out the connection with formal concept analysis.

REFERENCES

- Aerts, D. (1982). Description of many physical entities without the paradoxes encountered in quantum mechanics, *Found. Phys.* **12**, 1131–1170.
- Aerts, D. (1983). Classical theories and non classical theories as a special case of a more general theory, *J. Math. Phys.* **24**, 2441–2454.
- Aerts, D., Colebunders, E., Van der Voorde, A., and Van Steirteghem, B. (1999). State property systems and closure spaces: A study of categorical equivalence, *Int. J. Theor. Phys.* **38**, 359–385.
- Aumann, G. (1970). Kontakt-Relationen, *Sitz. Ber. Bayer. Akad. Wiss. Math. Nat. Kl.* **1970**, 67–77.
- Birkhoff, G. (1967). *Lattice Theory*, American Mathematical Society, Providence, Rhode Island.
- Erné, M. (1984). Lattice representations for categories of closure spaces, in *Categorical Topology*, Heldermann Verlag, Berlin, pp. 197–222.
- Ganter, B., and Wille, R. (1999). *Formal Concept Analysis*, Springer-Verlag, Berlin.
- Moore, D. J. (1995). Categories of representations of physical systems, *Helv. Phys. Acta* **68**, 658–678.

- Moore, D. J. (1999). On state spaces and property lattices, *Stud. Hist. Phil. Mod. Phys.*, **30**, 61–83.
- Moore, D. J. (2000). Fundamental structures in physics: A categorical approach, *Int. J. Theor. Phys.* this issue.
- Piron, C. (1976). *Foundations of Quantum Physics*, Benjamin, New York.
- Piron, C. (1989). Recent developments in quantum mechanics, *Helv. Phys. Acta* **62**, 82–90.
- Piron, C. (1990). *Mécanique quantique bases et applications*, Presses Polytechniques et Universitaires Romandes, Lausanne.